Acta Cryst. (1967). 23, 675 The Crystal Structure of the 2:1 Molecular Complex Between 1,3,7,9-Tetramethyluric Acid **and 3,4-Benzpyrene**

BY A. DAMIANI, E. GIGLIO AND A. M. LIQUORI

Centro Nazionale di Chimica delle Macromolecole (CNR) Sezione III, *Istituto Chimico,*

Universit3 di Napoli, Napoli, Italy

AND A. RIPAMONTI

Istituto di Chimica, Università di Trieste, Trieste, Italy

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Crystals of the 2:1 molecular complex between 1,3,7,9-tetramethyluric acid (TMU) and 3,4-benzpyrene (BP) are triclinic, space group P_1 , with $a = 9.33 \pm 0.02$, $b = 10.59 \pm 0.02$, $c = 10.84 \pm 0.02$ Å, $\alpha = 119^\circ 37'$ $\pm 12'$, $\beta = 113°20' \pm 12'$, $\gamma = 87°20' \pm 12'$, $Z = 1$. The crystal structure has been analysed by allowing systematic independent translations and rotations of the two TMU and the BP in their molecular planes and by calculating a large set of most significant structure factors which were compared with the experimental diffraction amplitudes. The structure is characterized by a plane to plane alternate stacking of two TMU and one BP molecule arranged in infinite columns parallel to the b axis. The molecular planes are almost exactly parallel with an average perpendicular separation of 3.48 Å. The strongest intermolecular forces are 'polarization bonding' and van der Waals interactions.

Introduction

As part of a research programme on the interactions between deoxyribonucleic acid and aromatic hydrocarbons (Liquori, Ascoli & Botrè, 1962; Liquori, De Lerma, Ascoli, Botrè & Trasciatti, 1962) the crystal structures of some molecular complexes between 1,3,7,9 tetramethyluric acid (TMU) and pyrene (PR), 3,4-benzpyrene (BP) and coronene (CR) have been undertaken. In previous papers the crystal structure of the 1 : 1 molecular complex between TMU and PR (Damiani, De Santis, Giglio, Liquori, Puliti & Ripamonti, 1965) and a preliminary account of the crystal structures of the 2:1 molecular complexes between TMU and BP and CR (Damiani, Giglio, Liquori, Puliti & Ripamonti, 1966, 1967) have been reported.

In this communication we wish to report the results of a three-dimensional crystal structure analysis carried out on the 2:1 molecular complex between TMU

Experimental

Prismatic crystals of the 2:1 TMU-BP complex were prepared from benzene according to Weil-Malherbe (1946) and used to collect the X-ray data. 740 reflexions, about 20% of the total accessible reflexions, were recorded with a Wiebenga integrating Weissenberg camera. The intensities of the reflexions *hOl, hll, . .., h5l, Okl, lkl, ..., 5kl* and those collected by recording along the $0\overline{1}\overline{1}$ diagonal were estimated photometrically, corrected for Lorentz and polarization factors and put on a common scale with the aid of the *hOl* intensity data. No absorption corrections were applied.

Crystal data

 $C_{20}H_{12}$. $2C_9N_4O_3H_{12}$, $M=700.7$. Yellow crystals, triclinic.

 $a=9.33 \pm 0.02$, $b=10.59 \pm 0.02$, $c=10.84 \pm 0.02$ Å, $\alpha = 119^{\circ}37' \pm 12'$, $\beta = 113^{\circ}20' \pm 12'$, $\gamma = 87^{\circ}20' \pm 12'$ $U=840 \text{ Å}^3$, $Z=1$, $F(0,0,0)=368$, λ Cu $K\alpha=1.542 \text{ Å}$. Space group P_1 . The theoretical density, 1.38 g.cm⁻³, is in good agreement with the measured one, 1.39 $g.cm^{-3}$, if one BP and two TMU molecules are in the unit cell.

Determination of the structure

By inspection of the three-dimensional weighted reciprocal lattice it was deduced that the TMU and BP molecules lie in almost parallel planes. The coefficients of the equations of the mean molecular planes were derived by means of very strong reflexions belonging to the levels with $k=3,6,9,12$. These planes intersect the y axis at approximately 0, $\frac{1}{3}b$ and $\frac{2}{3}b$.

The structure was determined by performing systematic independent translations and rotations of the two TMU and the BP in their molecular planes and

Table 1. *Comparison between observed and calculated structure factors* Values are multiplied by 10.

by computing the minimum residual of a set of most significant structure factors. During this trial process the calculated diffraction amplitudes were compared with the observed ones (Albano, Bellon, Pompa & Scatturin, 1963; Bhuiya & Stanley, 1963, 1964).

The analysis was started assuming for the BP a rigid model with bond lengths of 1.40 Å and bond angles of 120°, which are values not far away from the data found by Iball & Young (1956), whereas for the TMU Sutor's (1963) bond distances and angles were used. This procedure was repeated three times, using first 30, then 60 and finally 122 structure factors, corresponding to all the reflexions in the spheres of radius 2 sin $\theta/\lambda =$ 0.26 , 0.33 and 0.44 respectively. 15 minima were located for the reliability index R in the range $0.26-0.36$ for all 740 reflexions. Such a large number of minima is partly due, at low Bragg angles, to a pseudo twofold axis in the Fourier transform passing approximately, in real space, through $C(9)$ and $C(19)$ of BP and partly to a pseudo mirror plane in the TMU molecule passing through $C(8)$ and the middle point of the bond $C(2)$ - $N(1)$. Moreover two different TMU couplings, the first one resembling the centrosymmetrical arrangements found in crystals of pure TMU (Sutor, 1963; De Santis, Giglio & Liquori, 1960), the second one an iso-oriented arrangement, turned out to have scarcely any effect on the values of R .

At this point all the minima, except four belonging to the second coupling, were discarded because of some. short intermolecular distances. No criteria were found for preferring one rather than another of these four minima, which differ from each other mainly in the BP arrangement (Damiani, Giglio, Liquori & Ripamonti, 1967). The aromatic hydrocarbon takes in fact four orientations corresponding to a rotation through 180° in the molecular plane and to the above mentioned pseudo-twofold operation. Moreover it was clear that for these four arrangements even the refinement could not decide about the true structure.

Nevertheless the structure corresponding to the deepest minimum was refined somewhat by means of a Fourier differential synthesis computed to refine individual parameters of C , N and O atoms (Damiani, De Santis, Giglio & Ripamonti, 1963) as a test of the goodness of the analysis. The refinement was carried out with approximate anisotropic temperature factors, one for each chemical species (see Table 3). The disagreement index dropped from 0.26 to 0.18 , and Table 1 compares the observed and calculated structure factors, computed with the use of the atomic scattering factors of Berghuis, Haanappel, Potters, Loopstra, MacGillavry $&$ Veenendaal (1955) for C, N and O and that of McWeeny (1951) for H. The atomic parameters used in the structure factor calculations are reported in Table 2. In spite of the rather high final value of *and the low resolution, mainly caused by the small* percentage of the total accessible reflexions recorded, we feel that our results are an adequate representation of the complexity of the structure. For this purpose $C(20)$

Figs. 1, 2 and 3 show the good agreement between the three-dimensional Fourier syntheses observed and calculated, evaluated in the molecular planes of the two TMU and of the BP (Damiani, Giglio & Ripamonti, 1967).

Fig. 4 shows composite projections along the b axis of the electron density distributions calculated through sections containing the atoms of the first TMU, the

second TMU and the BP molecule. Fig. 5 shows a composite projection of the complex along the a axis. The difference between observed and calculated electron density is always lower than ± 0.5 e.Å⁻³.

At this stage eight atoms, six in BP and two in TMU molecules, are not resolved, as can be seen in the Fourier synthesis computed with observed data. It appeared clear that no further refinement would be profitable, mainly because of the low number of observed reflexions and the rather anomalous thermal motion of the atoms. The molecular packing of this crystalline structure is shown in Figs. 6, 7 and 8.

The similarity between the computed structure factors of the possible four structures selected as above described indicates that for all the models the same degree of uncertainty could be foreseen. A few refinement cycles on a second model confirmed it, and therefore no further refinement work has been carried out on the other models. It seemed therefore that the selection of a single model is not possible on the basis of the available experimental data.

The full description of the method used in the analysis and the rigid body coordinates corresponding to all the minima found in the R function is given in a separate paper (Damiani, Giglio, Liquori & Ripamonti, 1967).

All the calculations were performed on an IBM 1620 basic computer.

Fig. 1. Three-dimensional Fourier syntheses observed (a) and calculated (b) , computed in the molecular plane, of the first TMU. Contours are at intervals of $1e.\AA^{-3}$ starting with $2 e.A^{-3}.$

Fig.2. Three-dimensional Fourier syntheses observed (a) and calculated (b), computed in the molecular plane, of the second TMU. Contours are at intervals of 1 e. A^{-3} starting with $2 e.A^{-3}$.

Fig. 3. Three-dimensional Fourier syntheses observed (a) and calculated (b), computed in the molecular plane, of the BP. Contours lines 2, 3, 3.5, 4, 4.5, 5 e. A^{-3} .

Discussion

Because of the low resolution obtained in the crystal structure determination it is impossible to achieve detailed information about the molecular geometry. However the structure can be considered as established as far as the geometrical aspects common to all four models are concerned. In fact the other three possible models may be derived from this one through pseudosymmetry operations as previously described.

The least-squares planes of the BP and TMU molecules, calculated according to Schomaker, Waser, Marsh & Bergman (1959) with unitary weights for all the atoms, are:

The equations are referred to crystal axes and *x,y,z* are fractional coordinates. The molecular planes are nearly parallel with an average perpendicular interplanar spacing of 3.48 Å, and the structure is characterized by a plane to plane alternate stacking of two TMU and one BP molecule arranged sandwich-like in infinite columns parallel to the b axis. These geometrical features are very similar to those found in the TMU-PR complex and do not seem to indicate the existence of charge transfer forces, even if the projection of three molecules, TMU-BP-TMU, of the same stack on the BP plane (Fig. 6), shows a little more marked degree of overlap than in the pyrene complex.

Fig. 5. Composite drawing of the final electron density series, viewed along the a axis. Contours are at intervals of 1 e. A^{-3} starting with 1 e. \AA ⁻³.

Table 3. Temperature factors for the different atomic species in the form: $\exp[-(b_{11}h^2 + b_{12}hk + b_{13}hl + b_{22}k^2 + b_{23}kl + b_{33}l^2)]$

	D_{11}	v_{12}	D_{13}	b_{22}	023	b_{33}
С	0.0195	0.0045	0.0180	0.0120	0.0175	0.0190
N	0.0190	0.0045	0.0170	0.0115	0.0170	0.0185
о	0.0195	0.0045	0.0180	0.0120	0.0175	0.0190

Fig.6. Projection on to the BP least-squares plane showing the degree of overlap of an asymmetric unit and the packing with neighbouring molecules. The mean coordinate (A) of each molecule along the perpendicular to the plane of plotting is quoted. In a triplet the highest molecule is shown by heavy lines, the middle one by full light lines, and the lowest one by dashed lines.

Fig.7. Schematic drawing of the perpendicular separation of the molecules in the asymmetric unit.

Fig. 8. Clinographic view of the structure.

Spectroscopic properties of the complexes in solution and in the solid state (Van Duuren, 1964) support the view that the intermolecular attractions are due to 'polarization bonding' (Wallwork, 1961) and to van der Waals interactions. Moreover, it is interesting to note that the coupling of two adjacent TMU molecules, related by a pseudo glide plane, is compatible with an antiparallel direction of the TMU's π dipole moments which lie almost perpendicular to the long axis of the molecule.

Inspection of the three-dimensional Fourier syntheses suggests that the molecules are undergoing strong librations in their own planes as indicated by the elongation of the peripheral atoms. This is also supported by the temperature factors used in the molecular planes which are higher than those in the perpendicular direction. However, the presence in the X-ray photographs of some streaks, and the nearly identical orientations which either TMU or BP may adopt, indicate the possibility of disorder phenomena, not checked in this work, like those found in the indoles-trinitrobenzene complex (Hanson, 1964). In fact the molecular arrangements corresponding to the four minima fill the available space equally well and are practically equivalent with respect to the Fourier transform.

Like the least-squares refinement with rigid groups (Scheringer, 1963), the method of analysis used in this work can be considered as a very convenient one, particularly when the ratio between the experimental data and the parameters is not too high. Moreover, it must be emphasized that, in the minimum residual search, convergence can be reached even if few reflexions are used.

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Minimum Residual Analysis of Crystal Structures

By A. DAMIANI, E. GIGLIO AND A.M. LIQUORI

Centro Nazionale di Chimica delle Macromoleeole (CNR), Sezione III, *Istituto Chimico, Università di Napoli, Napoli, Italy*

AND A. RIPAMONTI

Istituto di Chimica, Università di Trieste, Trieste, Italy

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When the approximate geometrical features of the molecule or molecules present in a unit cell of a crystal are known it is possible to utilize a few X-ray experimental intensities of low Bragg angle to scan systematically both the orientation and the positions in the cell with the help of the 'minimum residual' function R. The regions in which a reasonable agreement has been observed can subsequently be explored more carefully, increasing the number of reflexions and decreasing the rotational and translational increments. Then computation of intermolecular distances and potential energy as well as rigid-body least-squares techniques can be powerfully coupled together in the search for the true minimum. In this way it is possible to locate the individual molecules in the unit cell with such an accuracy as to allow the subsequent use of standard refinement techniques. The result of such an analysis in the case of a trimolecular complex (2:1 complex of 1,3,7,9-tetramethyluric acid and 3,4-benzpyrene) is discussed in detail.

Introduction

If the geometry of a molecule is known its Fourier transform can be evaluated as originally done by Knott (1940) or reproduced through the 'optical diffractometer' derived by Taylor, Hinde & Lipson (1951) from the X-ray microscope of Bragg (1939). Since the first application by Lipson & Taylor (1951) its usefulness has been widely recognized and exploited, stimulating many investigations to enlarge the field of its applica-